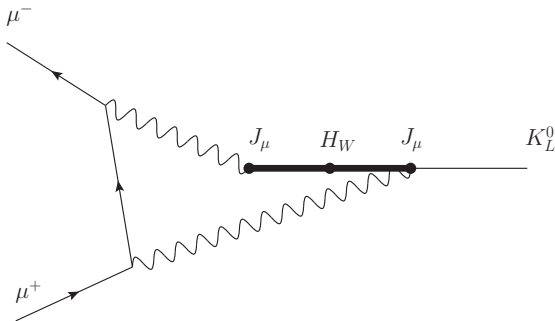


# Calculating the two-photon contribution to the real part of $\pi^0 \rightarrow e^+e^-$ decay amplitude

Norman Christ, Xu Feng, Luchang Jin, Cheng Tu, Yidi Zhao\*

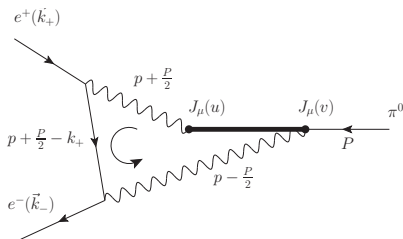
September 23, 2019

# Two-photon contribution to $K_L \rightarrow \mu^+ \mu^-$



- ▶ Intermediate states with lower energy than kaon mass:  
 $\gamma\gamma, |0\rangle, \eta, \pi, \pi\pi(\gamma)$ , etc.
- ▶ Calculation in Euclidean space would result in exponentially divergent behavior.

$$\pi^0 \rightarrow e^+ e^-$$



$$\langle e^+(k_+) e^-(k_-) | \pi^0(P) \rangle = \int d^4 u \, d^4 v \, H_{\mu\nu}(u, v) L^{\mu\nu}(u, v)$$

$$H_{\mu\nu}(u, v) = \langle 0 | J_\mu(u) J_\nu(v) | \pi \rangle$$

$$L^{\mu\nu}(u, v) = \int d^4 p \, e^{-ip \cdot (u-v)} \left[ \frac{g_{\mu\mu'}}{(p + \frac{P}{2})^2 - i\epsilon} \right] \left[ \frac{g_{\nu\nu'}}{(p - \frac{P}{2})^2 - i\epsilon} \right] \\ \bar{u}(k_-) \gamma_{\mu'} \left[ \frac{\gamma \cdot (p + \frac{P}{2} - k_-) + m_e}{(p + \frac{P}{2} - k_-)^2 + m_e^2 - i\epsilon} \right] \gamma_{\nu'} v(k_+)$$

$$\pi^0 \rightarrow e^+ e^-$$

Let  $w$  be the relative coordinate of two EM currents i.e.  
 $w = u - v$ . Decay amplitude:

$$\begin{aligned} \mathcal{A} = & \int d^4 w \langle 0 | T \{ J_\mu(\frac{w}{2}) J_\nu(-\frac{w}{2}) \} | \pi^0 \rangle \\ & \int d^4 p e^{-ip \cdot w} \left[ \frac{g_{\mu\mu'}}{(p + \frac{P}{2})^2 + m_\gamma^2 - i\epsilon} \right] \left[ \frac{g_{\nu\nu'}}{(p - \frac{P}{2})^2 + m_\gamma^2 - i\epsilon} \right] \\ & \bar{u}(k_-) \gamma_{\mu'} \left[ \frac{\gamma \cdot (p + \frac{P}{2} - k_-) + m_e}{(p + \frac{P}{2} - k_-)^2 + m_e^2 - i\epsilon} \right] \gamma_{\nu'} v(k_+) \end{aligned} \quad (1)$$

Here the matrix element is a Minkowski-space quantity.

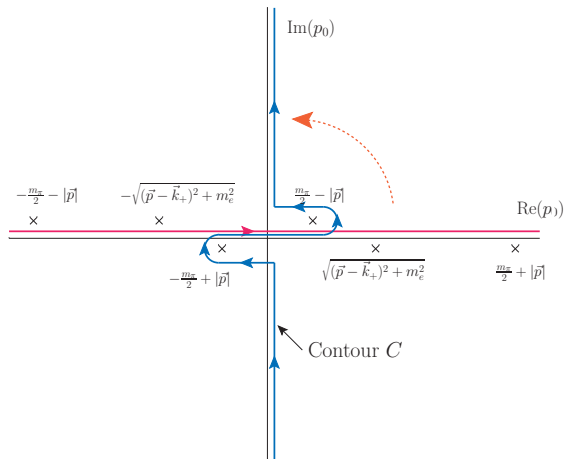
To calculate it using Lattice QCD, it is necessary to do a Wick rotation  $w^0 \rightarrow iw^0$ .

# Wick Rotation

$$L_{\mu\nu}(w) = \int d^4p \, e^{-ip \cdot w} \left[ \frac{g_{\mu\mu'}}{(p + \frac{P}{2})^2 - i\epsilon} \right] \left[ \frac{g_{\nu\nu'}}{(p - \frac{P}{2})^2 - i\epsilon} \right] \bar{u}(k_-) \gamma_{\mu'} \left[ \frac{\gamma \cdot (p + \frac{P}{2} - k_-) + m_e}{(p + \frac{P}{2} - k_-)^2 + m_e^2 - i\epsilon} \right] \gamma_{\nu'} v(k_+)$$

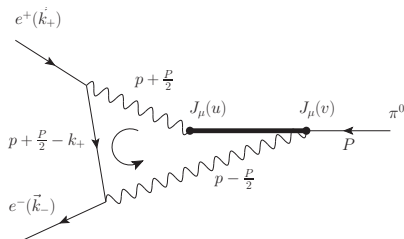
- ▶ After  $w^0$  is Wick-rotated, the exponential in  $L_{\mu\nu}(w)$  introduces exponential growth.
- ▶ The  $p^0$  contour also needs to be rotated.
- ▶ Because of the presence of intermediate states with lower energy than the mass of pion, a naive Wick rotation  $p^0 \rightarrow ip_E^0$  will make the integral blow up.

# Wick Rotation



- When  $|\vec{p}| < \frac{m_\pi}{2}$ , the  $p^0$  contour needs to be deformed to circumvent the two poles that cross the imaginary axis.

# Wick Rotation

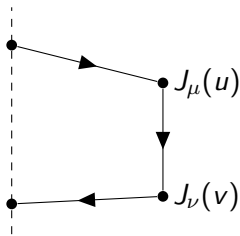


- ▶  $L_{\mu\nu} \propto e^{M_\pi \frac{|w^0|}{2}}$ , where  $w^0 = u^0 - v^0$
- ▶ The lightest intermediate state between  $J_\mu$  and  $J_\nu$  is  $\pi\pi$  state  
 $\rightarrow H_{\mu\nu} = \langle 0 | J_\mu(\frac{w}{2}) J_\nu(-\frac{w}{2}) | \pi \rangle \propto e^{-(E_n - \frac{M_\pi}{2})|w^0|}$ ,  $E_n > 2M_\pi$
- ▶ Amplitude converges exponentially at least as fast as  $e^{-M_\pi |w^0|}$ . We can safely perform the Euclidean space calculation.

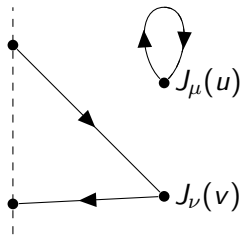
## Hadronic part

$$\langle 0 | T J_\mu(0) J_\nu(x) | \pi \rangle = \lim_{t \rightarrow -\infty} \frac{2m_\pi}{N_\pi} Z_V^2 e^{m_\pi |t|} \langle 0 | T J_\mu(0) J_\nu(x) \pi(t) | 0 \rangle$$

where  $N_\pi$  is the pion ground state amplitude  $N_\pi = \langle \pi | \pi(0) | 0 \rangle$ ,  $Z_V$  is the coefficient of EM current that connects the local non-conserved current with global conserved current.



(a) Connected Diagram



(b) Disconnected Diagram

- Contribution from the second graph is small (suppressed by  $SU(3)$  flavor symmetry).



# Summary

- ▶ Leptonic part integral is calculated numerically with CUBA library.
- ▶ Hadronic part can be extracted from three point function.

$$\langle 0 | T J_\mu(0) J_\nu(x) | \pi \rangle = \lim_{t \rightarrow -\infty} \frac{2m_\pi}{N_\pi} Z_V^2 e^{m_\pi |t|} \langle 0 | T J_\mu(0) J_\nu(x) \pi(t) | 0 \rangle$$

What we have calculated:

- ▶ Imaginary part of the decay amplitude  
It can be compared with Optical theorem prediction to help us estimate error in lattice calculation.
- ▶ Real part of the decay amplitude (Our real goal)

# Ensembles

## 1. $24^3 \times 64$ Iwasaki + DSDR

Lattice volume:  $24^3 \times 64$

$$a^{-1} = 1 \text{ GeV}$$

$$am_\pi = 0.13975(10)$$

Number of trajectories: 35

## 3. $32^3 \times 64$ Iwasaki + DSDR

Lattice volume:  $32^3 \times 64$

$$a^{-1} = 1.37 \text{ GeV}$$

$$am_\pi = 0.10468(32)$$

Number of trajectories: 17

## 2. $32^3 \times 64$ Iwasaki + DSDR

Lattice volume:  $32^3 \times 64$

$$a^{-1} = 1 \text{ GeV}$$

$$am_\pi = 0.139474(96)$$

Number of trajectories: 40

## 4. $48^3 \times 96$ Iwasaki

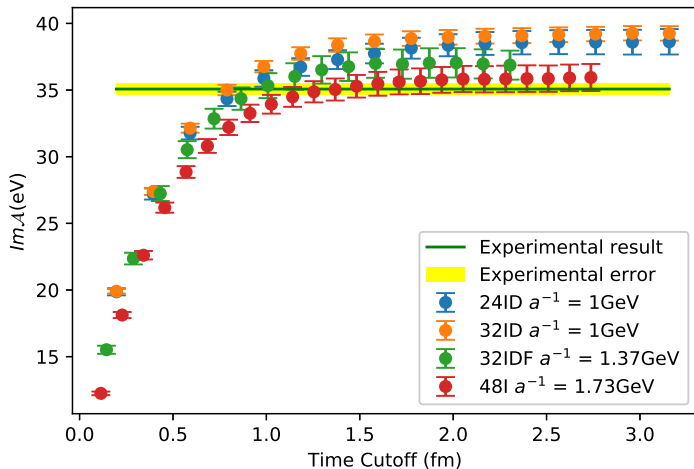
Lattice volume:  $32^3 \times 64$

$$a^{-1} = 1.73 \text{ GeV}$$

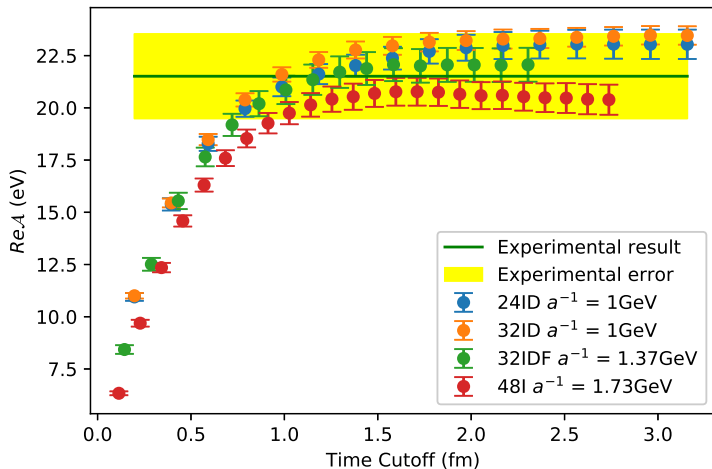
$$am_\pi = 0.08049(13)$$

Number of trajectories: 20

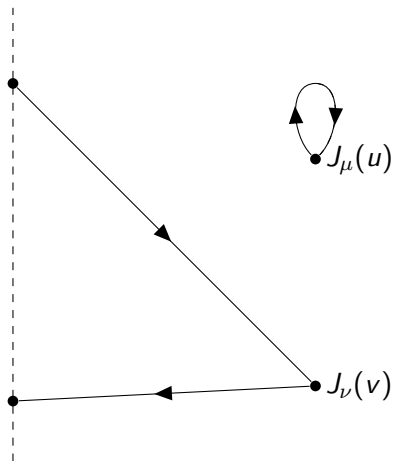
# Imaginary Part Contribution to Decay Rate



# Real Part Contribution to Decay Rate



# Disconnected Diagram



# Disconnected Diagram

- ▶ EM current tadpole  $Tr[\gamma_\mu S(x, x)]$  is calculated using random grid source.
- ▶ Disconnected diagram is about 2% of the connected diagram and they have the opposite sign.
- ▶ Disconnected diagram is treated as a part of systematic error.

# Conclusion

Experimental result from PDG:

$$\mathrm{Im}\mathcal{A} = 35.07(37) \text{ eV} \quad (2)$$

$$\mathrm{Re}\mathcal{A} = 21.51(2.02) \text{ eV} \quad (3)$$

---

Dispersion relation [Weil et al., 2017]:

$$\mathrm{Re}\mathcal{A} = 20.16(23) \text{ eV}$$

The error mostly comes from the experimental pion life time.

---

The result we obtained from lattice calculation:

$$\mathrm{Im}\mathcal{A} = 35.94(1.01)(1.19) \text{ eV} \quad (4)$$

$$\mathrm{Re}\mathcal{A} = 20.39(72)(70) \text{ eV} \quad (5)$$



# Conclusion

- ▶ We developed a method for dealing with two-photon intermediate state and combining QED and QCD part of the amplitude
- ▶ Using this method, we carried out a first-principles calculation of  $\pi^0 \rightarrow e^+e^-$  decay
- ▶ The next step is to calculate a more interesting process – the two photon contribution to  $K_L \rightarrow \mu^+\mu^-$  decay